

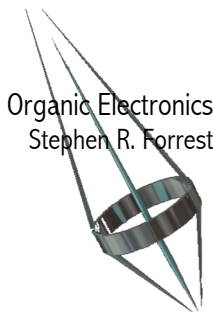
Week 2-6

Optical Detectors 1

Photodetection Basics

Organic photoconductors and photodiodes

Chapter 7.1-7.2



Organic Electronics
Stephen R. Forrest

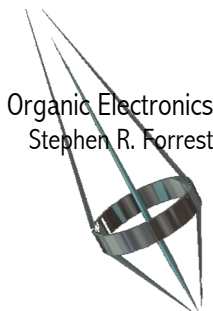
Objectives

- Understand the physics of photodetection in organic photoconductors and photodiodes
- Understand OPD performance characteristics
 - Dark current
 - Efficiency and responsivity
 - Bandwidth
 - Noise
- Learn about OPD applications
- Solar cells: what makes OPVs a compelling story?
- Learn how to characterize solar cell performance
- Solar cell architectures
 - Thermodynamic efficiency limits to single junction cells
 - Multijunction cells and other architectures
 - The role of morphology
 - Some materials
- What lies beyond the horizon?



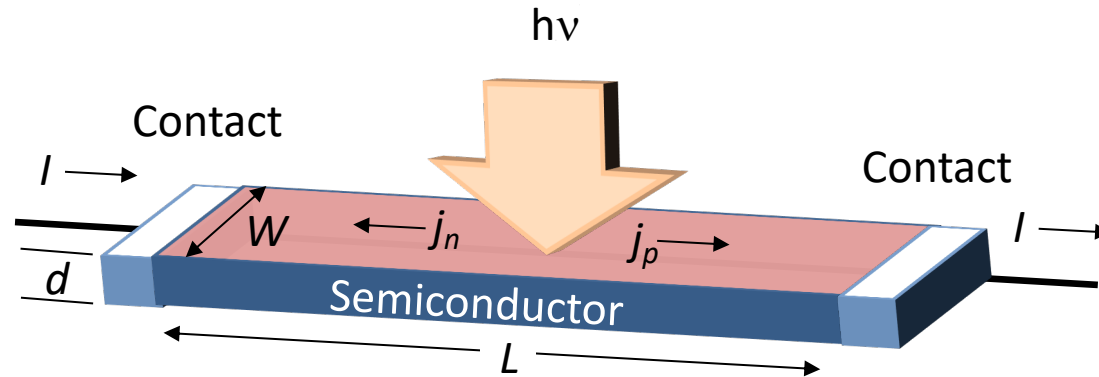
Photodetectors

- Transducers that convert light to another energy form (in our case, electricity)
- Types
 - Photoconductors
 - Photodiodes
 - These are operated in the reverse-biased (photodetection) or photovoltaic mode
- Properties
 - Sensitivity & Efficiency
 - Spectral range
 - Bandwidth
 - Dynamic range



Photoconductors

- Earliest organic electronic devices
- Simplest (no HJs needed)



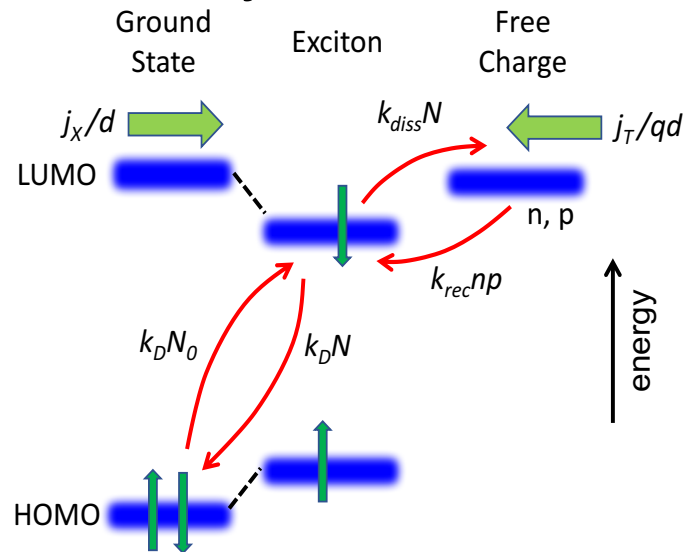
When illuminated, conductivity changes

$$\sigma = q(\mu_n n + \mu_p p) \left\{ \begin{array}{l} p = p_{ph} + p_0 \\ n = n_{ph} + n_0 \end{array} \right. \quad \boxed{n_{ph} = p_{ph}}$$

Without background doping: $n_0 = p_0 = n_i$

Photocharge generation

- Generation does not occur through an intermediate CT state as it does at OPD heterojunctions:



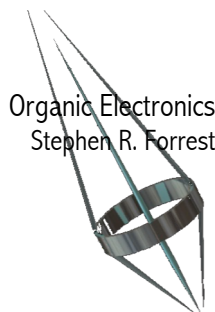
Generation rate: $G_{ph} = k_D n_{ph} = \frac{\eta_{ext} (P_{inc} \lambda / hc)}{dWL}$

$\tau_D = 1/k_D =$ lifetime of charge

$\eta_{ext} =$ external quantum efficiency (electrons out/photons in)

\Rightarrow Photocurrent:

$$j_{ph} = \sigma F = q n_{ph} (\mu_n + \mu_p) \frac{V_a}{L} = q \frac{\eta_{ext} (P_{inc} \lambda / hc)}{k_D} (\mu_n + \mu_p) \frac{V_a}{dWL^2}$$



Quantum Efficiency and Responsivity

$$\text{External quantum efficiency} = \frac{\text{No. electrons generated}}{\text{No. of photons incident}} = \eta_{ext} = \frac{hcj_{ph}}{q\lambda P_{inc}}$$

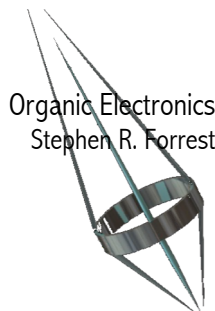
$$\text{Internal quantum efficiency} = \frac{\text{No. electrons generated}}{\text{No. of photons absorbed}} = \eta_{int} = \frac{hcj_{ph}}{q\lambda n_{ph} P_{inc}}$$

$$\text{where: } n_{ph} = \frac{(1-R)}{d} \int_0^d \exp(-\alpha(\lambda)x) dx$$

for a total reflection coeff't, R , from the surface, and an absorption coeff't of α in an active region of thickness, d .

Note: the total thickness must account for internal reflections and other cavity effects

$$\text{Responsivity} = \frac{\text{Current generated}}{\text{Power incident}} = \mathcal{R} = \frac{j_{ph}}{P_{inc}} = \frac{q\lambda}{hc} \eta_{ext} \quad [\text{A/W}]$$



Gain and bandwidth

$$j_{ph} = \sigma F = qn_{ph}(\mu_n + \mu_p)\frac{V_a}{L} = q\frac{\eta_{ext}(P_{inc}\lambda/hc)}{k_D}(\mu_n + \mu_p)\frac{V_a}{dWL^2}$$

⇒ A photoconductor has gain: $g = \frac{j_{ph}}{j_0} = \tau_D(\mu_n + \mu_p)\frac{V_a}{L^2}$

Where: $j_0 = q\eta_{ext}(P_{inc}\lambda/hc)/dW$

That is: $\text{gain} = \tau_D / t_{tr}$, where the carrier transit time is $t_{tr} = L/v = L/\mu F = L^2/\mu V$

$$g\eta_{ext} = \frac{j_{ph}A}{q(P_{inc}\lambda/hc)}$$

Quantum efficiency cannot be separated from gain

- Bandwidth: $\Delta f = 1/2\pi\tau_D$
- Leading to a **gain-bandwidth product**: $g\Delta f = 1/2\pi t_{tr}$



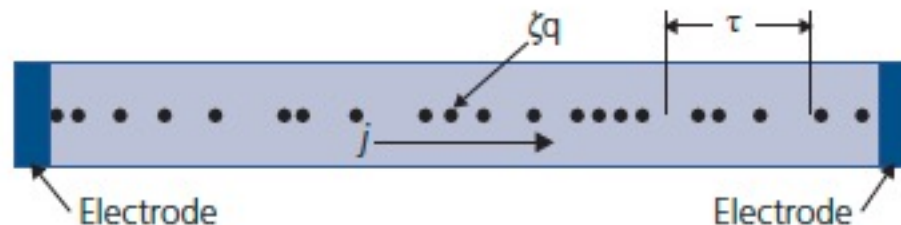
Calculating the Noise Current

(after Rose, 1963. *Concepts in Photoconductivity and Allied Problems*)

- Determines the sensitivity of a photodetector to low intensity signals

- Signal-to-noise ratio: $\frac{S}{N} = \frac{\langle i_{ph}^2 \rangle}{\langle i_n^2 \rangle} > \textcircled{1}$ $\langle i_{ph}^2 \rangle =$ mean square photocurrent
 $\langle i_n^2 \rangle =$ mean square noise current
minimum level of detectability

Consider a “general” photodetector. It has randomly generated particles, each carrying charge ζq in time interval, τ , between electrodes, resulting in current, j .



Then, the noise current is: $\langle i_n^2 \rangle^{1/2} = \frac{\langle n \rangle^{1/2}}{\tau} \zeta q$

where $\langle n \rangle^{1/2}$ is the rms number of particles collected in τ .

Calculating Noise Current, con't

Thus, in terms of the total mean current, i_T , the mean square noise current is:

$$\langle i_n^2 \rangle = \frac{\langle n \rangle}{\tau^2} (\zeta q)^2 = \frac{q i_T \zeta}{\tau}$$

Since the bandwidth is $\Delta f = 1/2\tau$, and accounting for both generation and recombination, we get a **shot noise current** of:

$$\langle i_s^2 \rangle = 4qgi_T\Delta f$$

If diffusion is dominant, then the charge delivered per particle is reduced by the fraction of charge diffusing to the contacts for a slab of length, L : $\zeta = L_D/L$.

Using $L_D = \sqrt{D\tau}$ and the Einstein relation for mobility, we obtain the **thermal**, or **Johnson noise**:

$$\langle i_{th}^2 \rangle = \frac{4k_B T \Delta f}{R_{PC}}$$

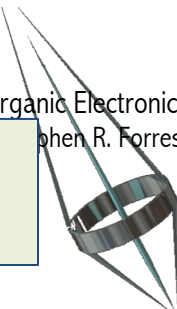
R_{PC} is the resistance of the conductor

Finally, there is **flicker**, or **1/f noise**:

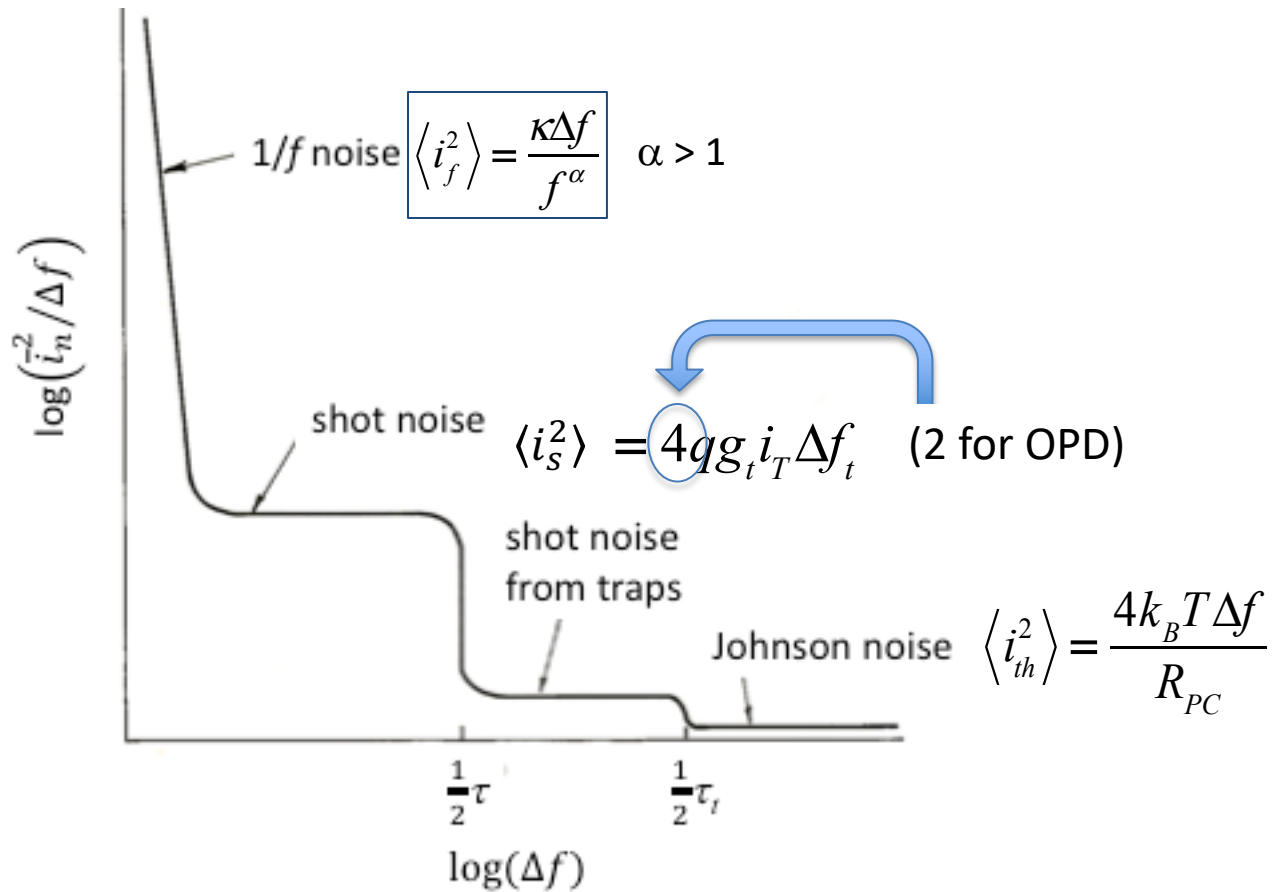
$$\langle i_f^2 \rangle = \frac{\kappa \Delta f}{f^\alpha}$$

κ , α are empirical constants

The total noise current is then the sum of the squares of the various contributions (they are uncorrelated): $\langle i_n^2 \rangle = \langle i_s^2 \rangle + \langle i_{th}^2 \rangle + \langle i_f^2 \rangle + \dots$

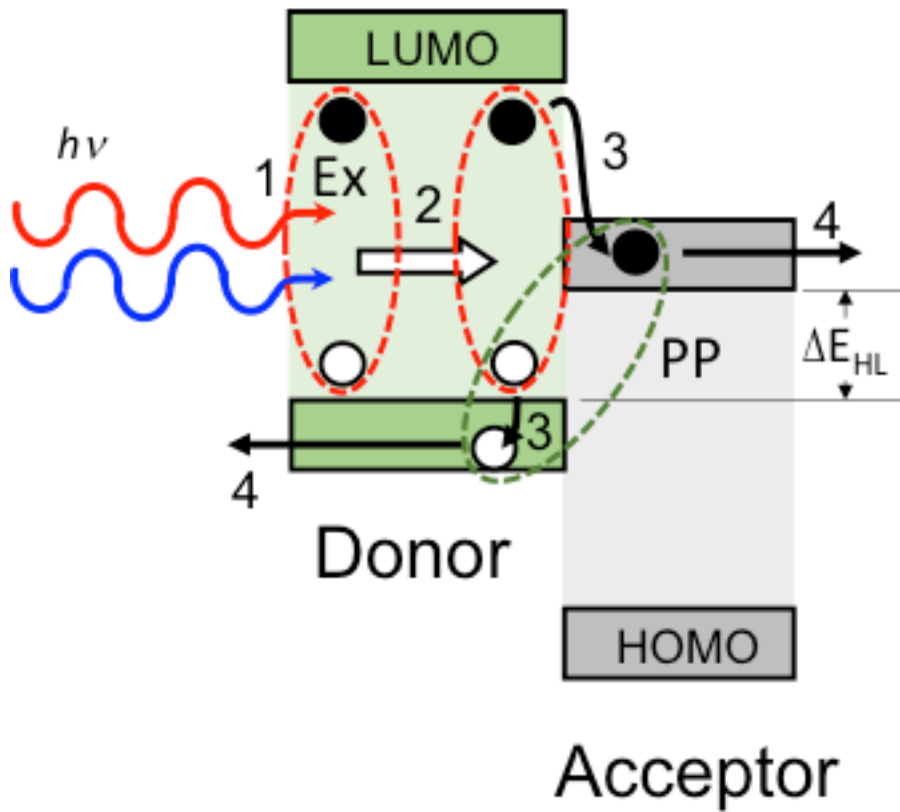


Graphically, Noise Spectra Look Like...



Photodiodes and solar cells

- Many of the same considerations as photoconductors except there is a junction for efficient charge separation.



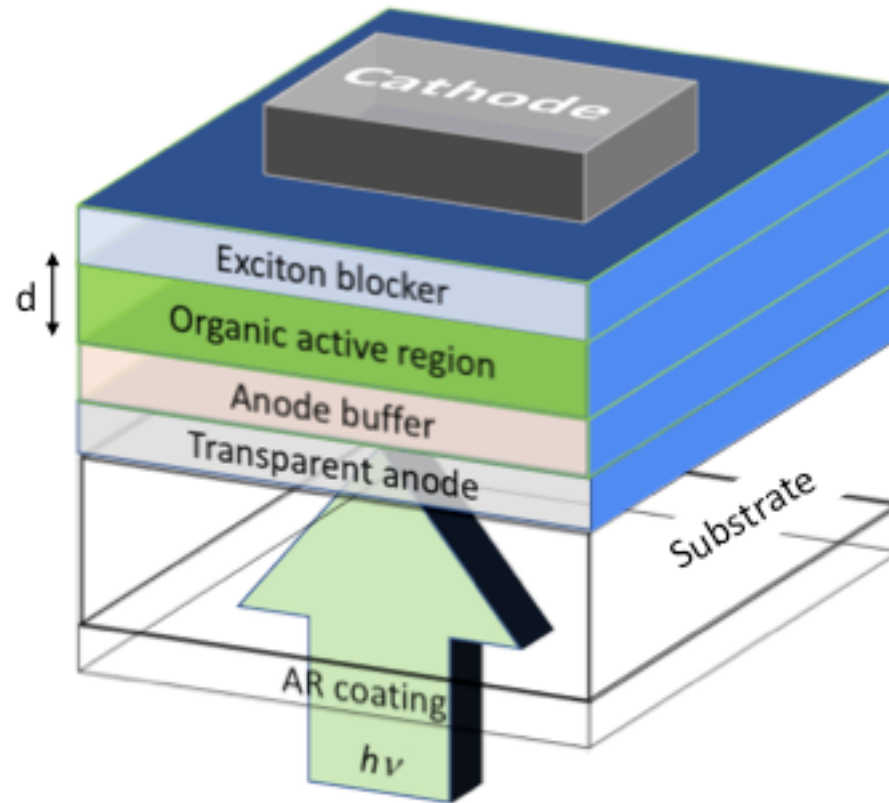
- Exciton generation by absorption of light (abs length $\sim 1/\alpha$)
- Exciton diffusion over $\sim L_D$
- Exciton dissociation by rapid and efficient charge transfer
- Charge extraction by the internal electric field

Typically: $L_D \ll 1/\alpha$

$$\eta_{ext} = \eta_A \eta_{ED} \eta_{CT} \eta_{CC}$$

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Basic OPD/OPV structure

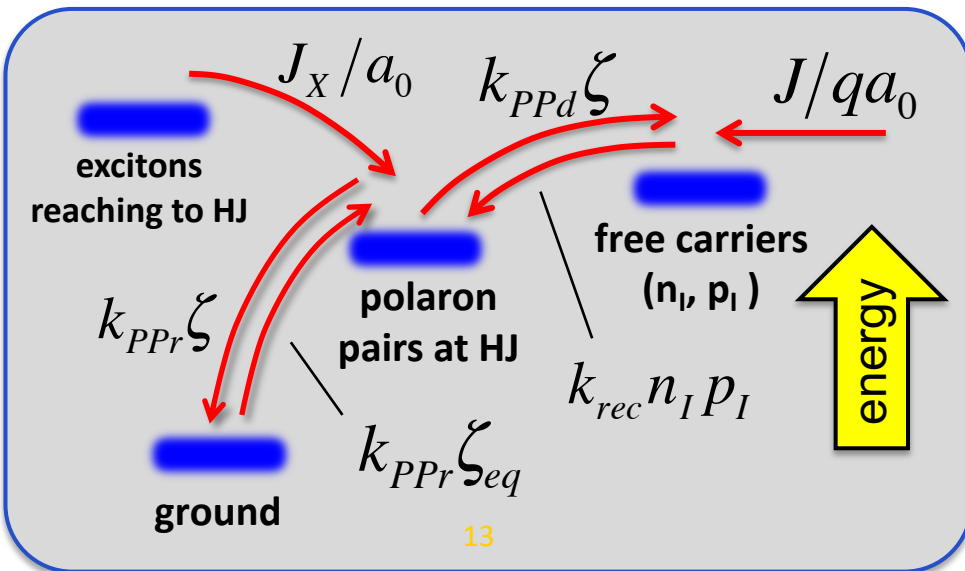


Current generation

- Recall (Ch. 4) that the j - V characteristics are given by:

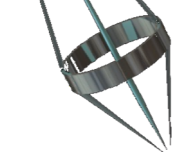
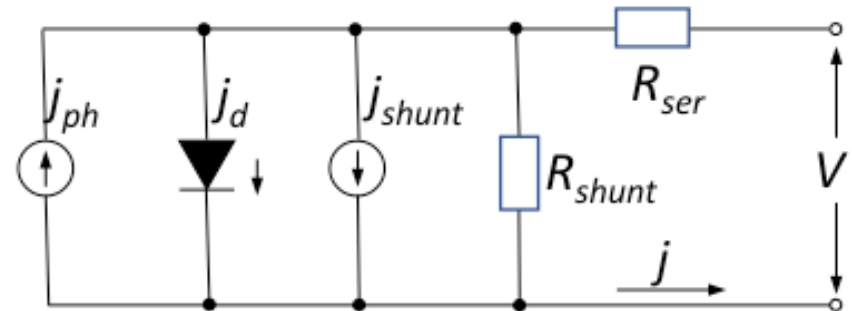
$$j = j_0 \left[\exp\left(\frac{q(V_a - jAR_{ser})}{n_S k_B T}\right) - \frac{k_{PPd}}{k_{PPd,eq}} \right] + \frac{V_a - jAR_{ser}}{R_{shunt}} - j_{ph}$$

Saturation current $j_0 = qa_0 k_{rec} N_S^2 (1 - \eta_{PPd}) \exp(-\Delta E_{HL} / k_B T)$

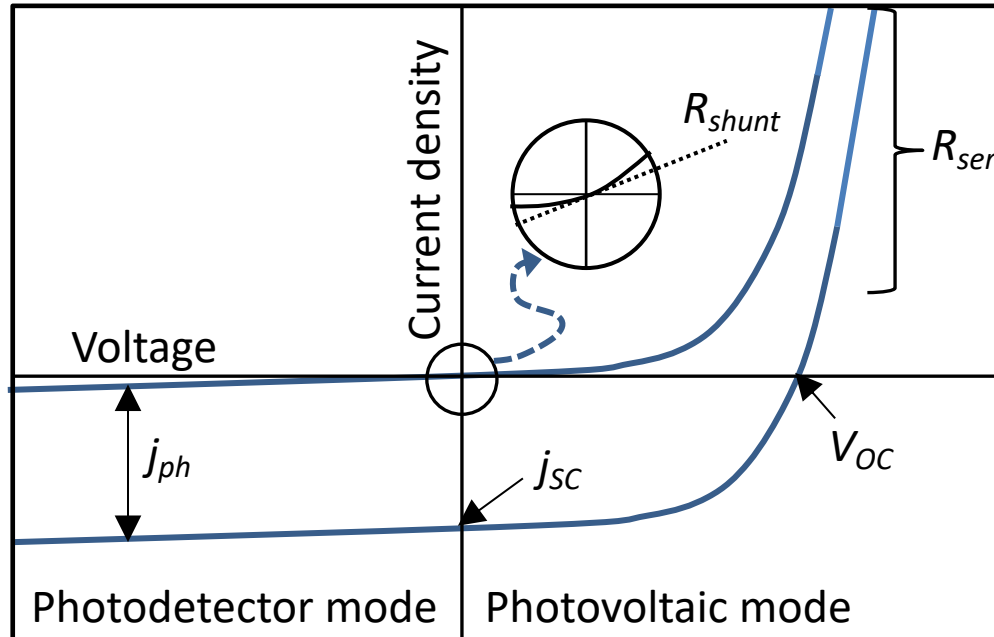


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Equivalent circuit

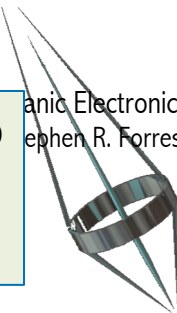


Current-Voltage Characteristics

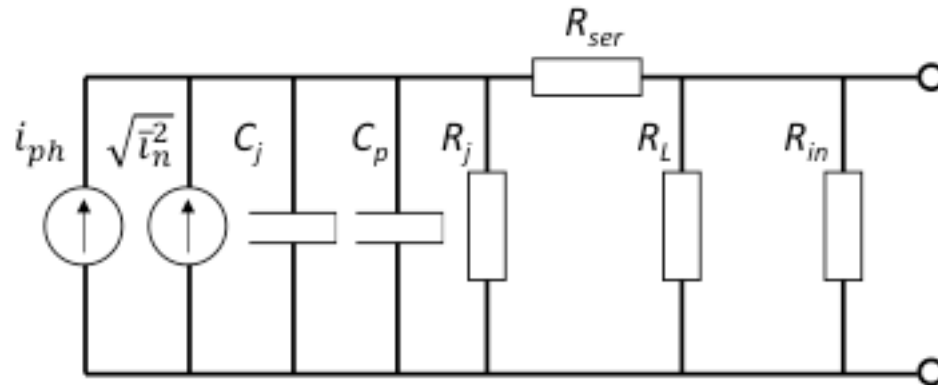


$$R_{shunt} = \left. \frac{1}{A} \frac{dV_a}{dj} \right|_{V_a=0}$$

- In the photovoltaic mode, the power is $P = jV < 0$; i.e. the device delivers power to the external circuit.
- In the photodetector mode, $P > 0$ and the detector dissipates power.



Photodetector Equivalent Circuit & Frequency Response



$$\Delta f = \frac{1}{2\pi} \left(\frac{1}{t_{tr}} + \frac{1}{\tau_{ED}} + \frac{1}{\tau_{RC}} \right)$$

$$\tau_{RC} = (R_{ser} + R_L || R_{in})(C_j + C_p) \quad (R_j \rightarrow \infty) \quad : \text{RC time constant}$$

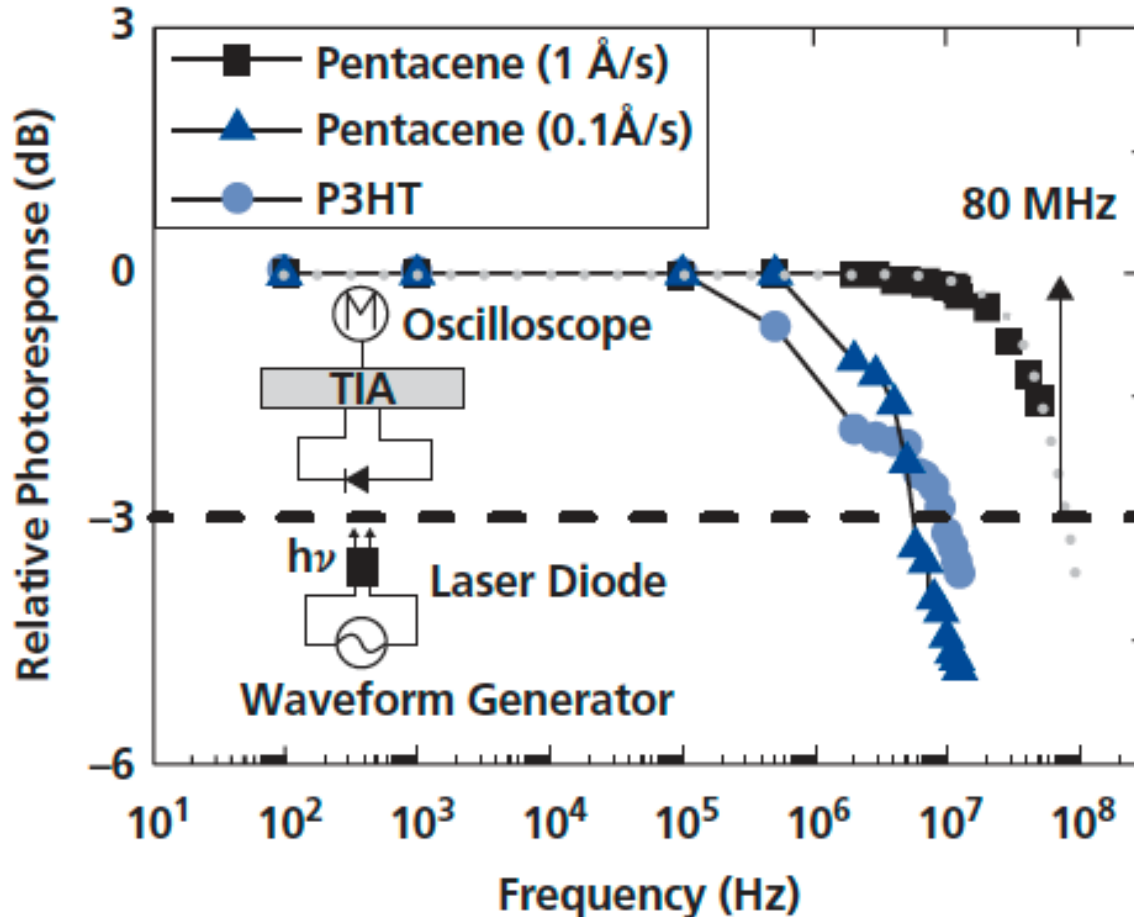
τ_{ED} : exciton diffusion time

$t_{tr} = w^2 / \mu V$: transit time through depleted regions of the device (w)

Gain-Bandwidth product = $g\Delta f = \Delta f$ since in a PD, $g = 1$.

Pentacene/C₆₀ OPD Frequency Response

High frequency response due to high pentacene mobility



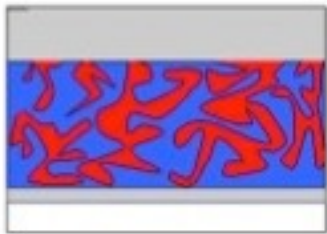
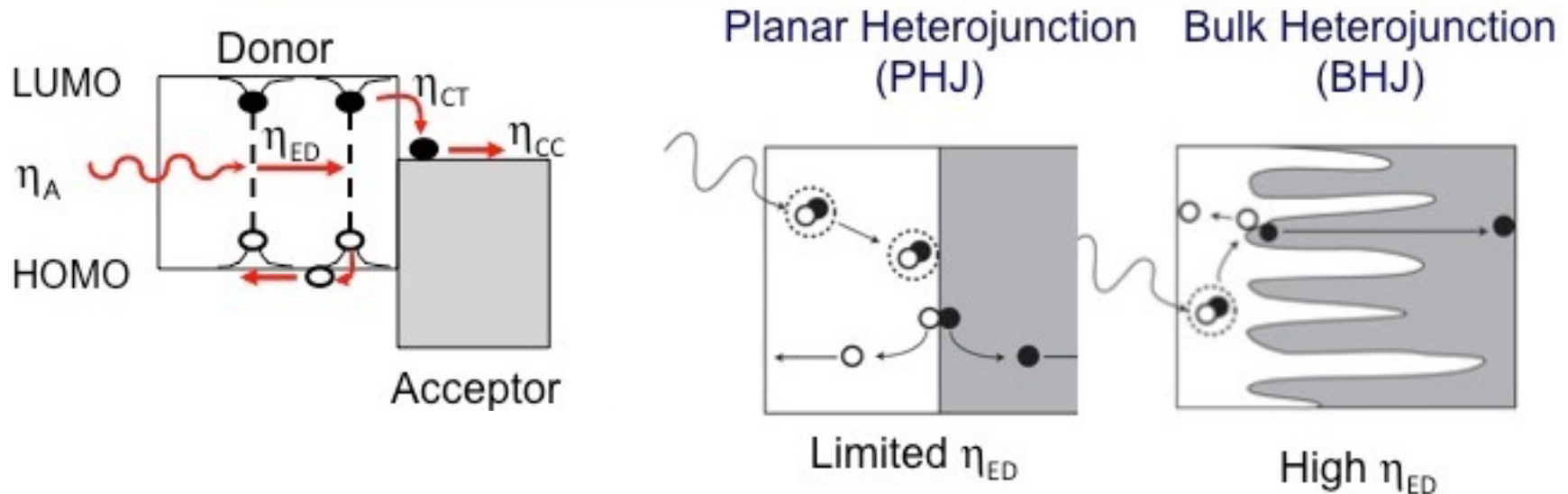
TIA: Transimpedance amplifier through which the diode is biased

Tsai et al. Appl. Phys. Lett., 95, 213308 (2009)

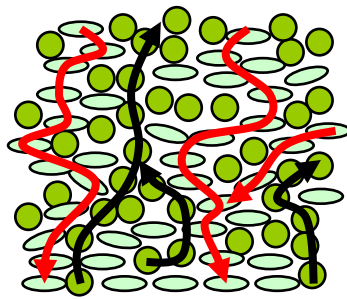


Heterojunction Morphologies

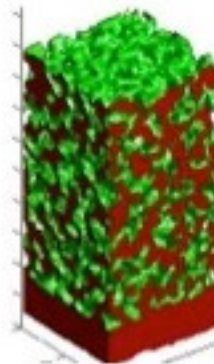
Breaking the tradeoff between L_D and α with BHJs



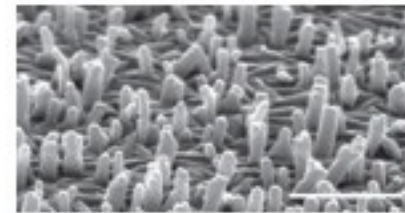
Bulk HJ



Mixed HJ

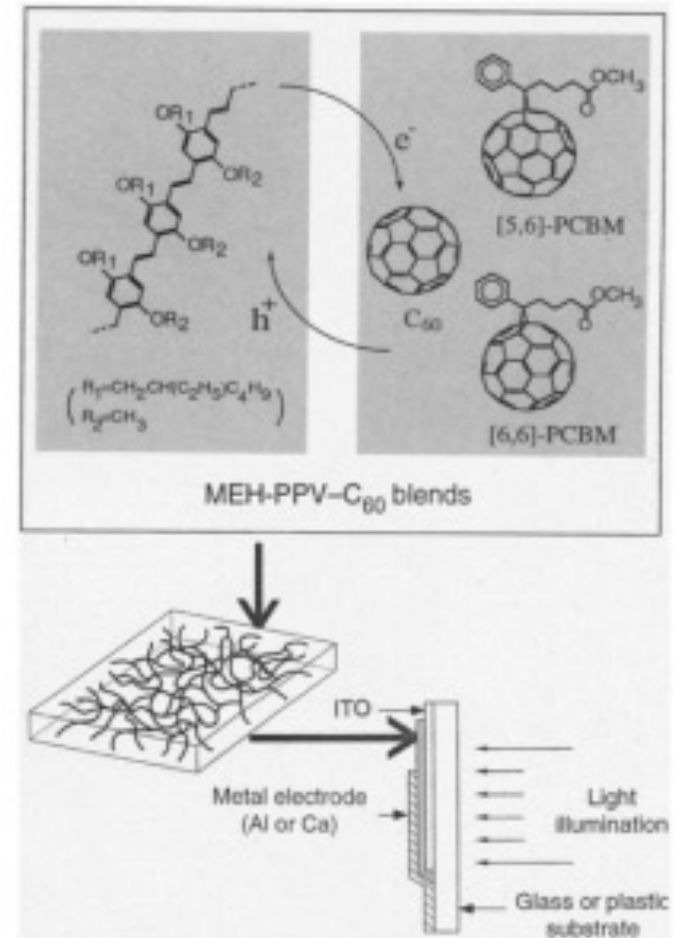
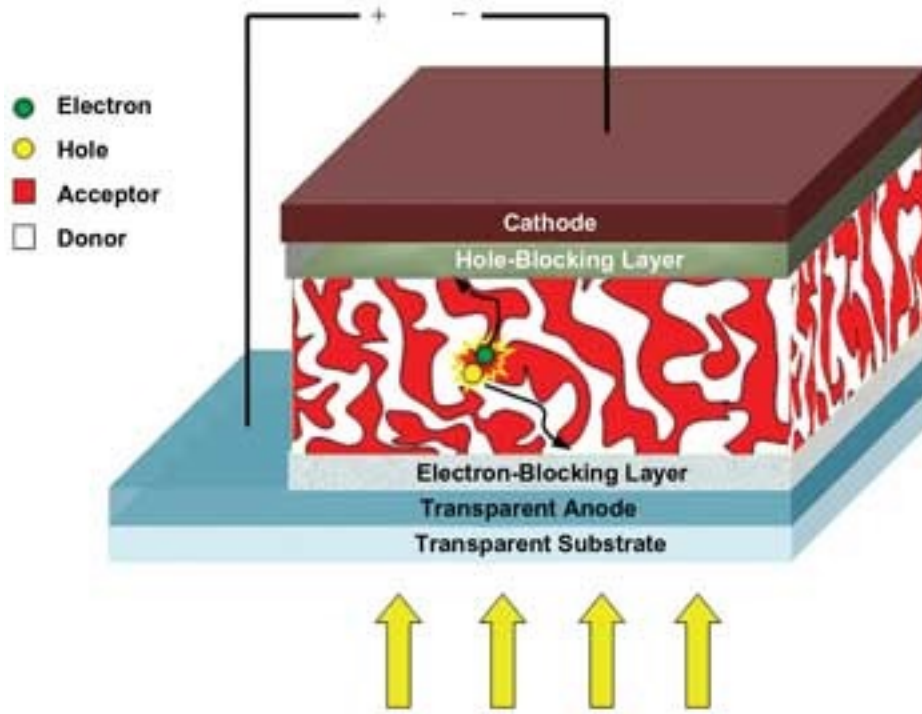


Annealed BHJ



Controlled BHJ

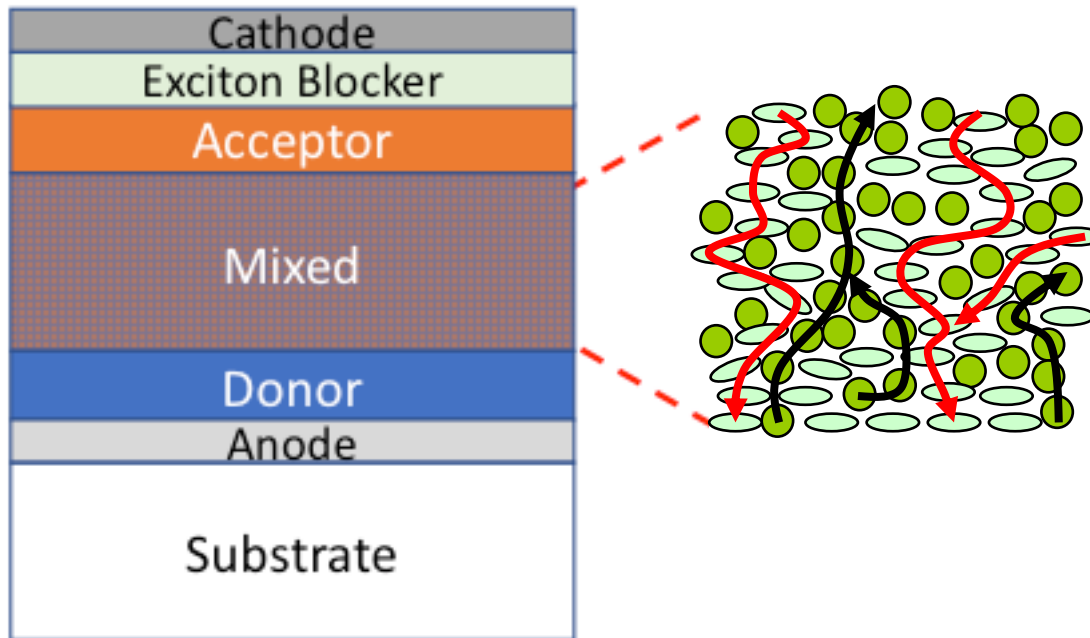
Polymer Bulk HJ



Yu et al. Science, **270**, (1995), 1789
 Halls et al., (1995) Nature, 376, 498.

Small Molecule Planar-Mixed HJ

Small molecule blends: $\eta_{ED} = 1$



$$\eta_{CC} = \frac{L_C}{x_M} (1 - \exp(-x_M/L_C))$$

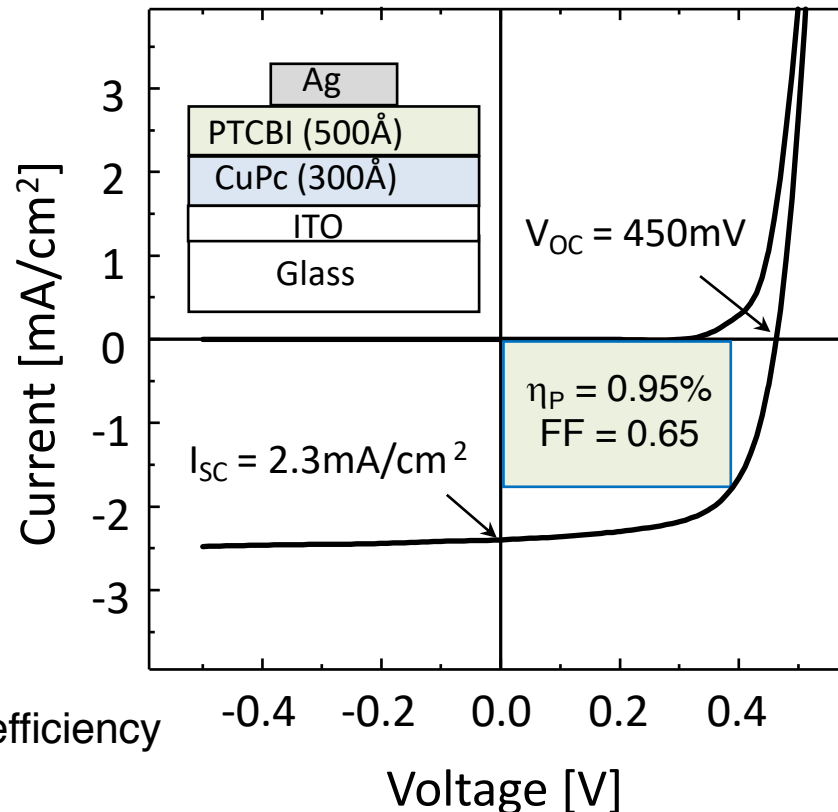
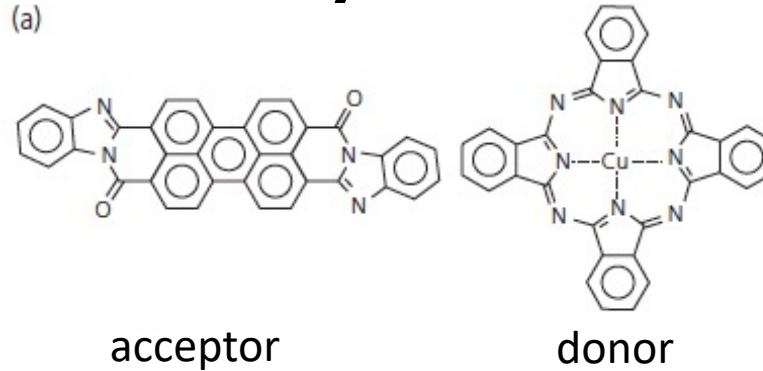
Charge carrier collection length, L_C , replaces diffusion length since excitons dissociate at point of generation without diffusion to HJ



Comparison of OPDs and OPCs

| Parameter | Photoconductor | Photodiode |
|--|---|---|
| Operating voltage | Near equilibrium ($V_a \rightarrow 0$) | Reverse bias |
| Photocurrent gain (g) | τ/t_{tr} (1-10 ⁶) | 1 |
| η_{int} | $k_{diss} / (k_{diss} + k_D)$ | $k_{ppd} / (k_{PPd} + k_{PPr})$ |
| η_{ext} | $\frac{j_{ph} A}{qg(P_{inc} \lambda/hc)}$ | $\frac{j_{ph} A}{q(P_{inc} \lambda/hc)}$ |
| Responsivity | $qg\eta_{ext} (\lambda/hc)$ | $q\eta_{ext} (\lambda/hc)$ |
| Bandwidth (Δf) | $1/2\pi\tau_D$ | $1/2\pi t_{tr}$ |
| Gain-bandwidth product ($g\Delta f$) | $1/2\pi t_{tr}$ | $1/2\pi t_{tr}$ |
| $\bar{i}_n^2 / \Delta f$ | $(4k_B T) / R_{PC} + \kappa / f^\alpha$ | $2qi_T + 4k_B T / R_L \parallel R_{in}$ |
| Specific detectivity (D^*) | $q\eta_{ext} (\lambda/hc) \sqrt{\frac{A}{(4k_B T) / R_{PC} + \kappa / f^\alpha}}$ | $q\eta_{ext} (\lambda/hc) \sqrt{\frac{A}{2qi_T + 4k_B T / R_L \parallel R_{in}}}$ |

The first bilayer OPD/OPV



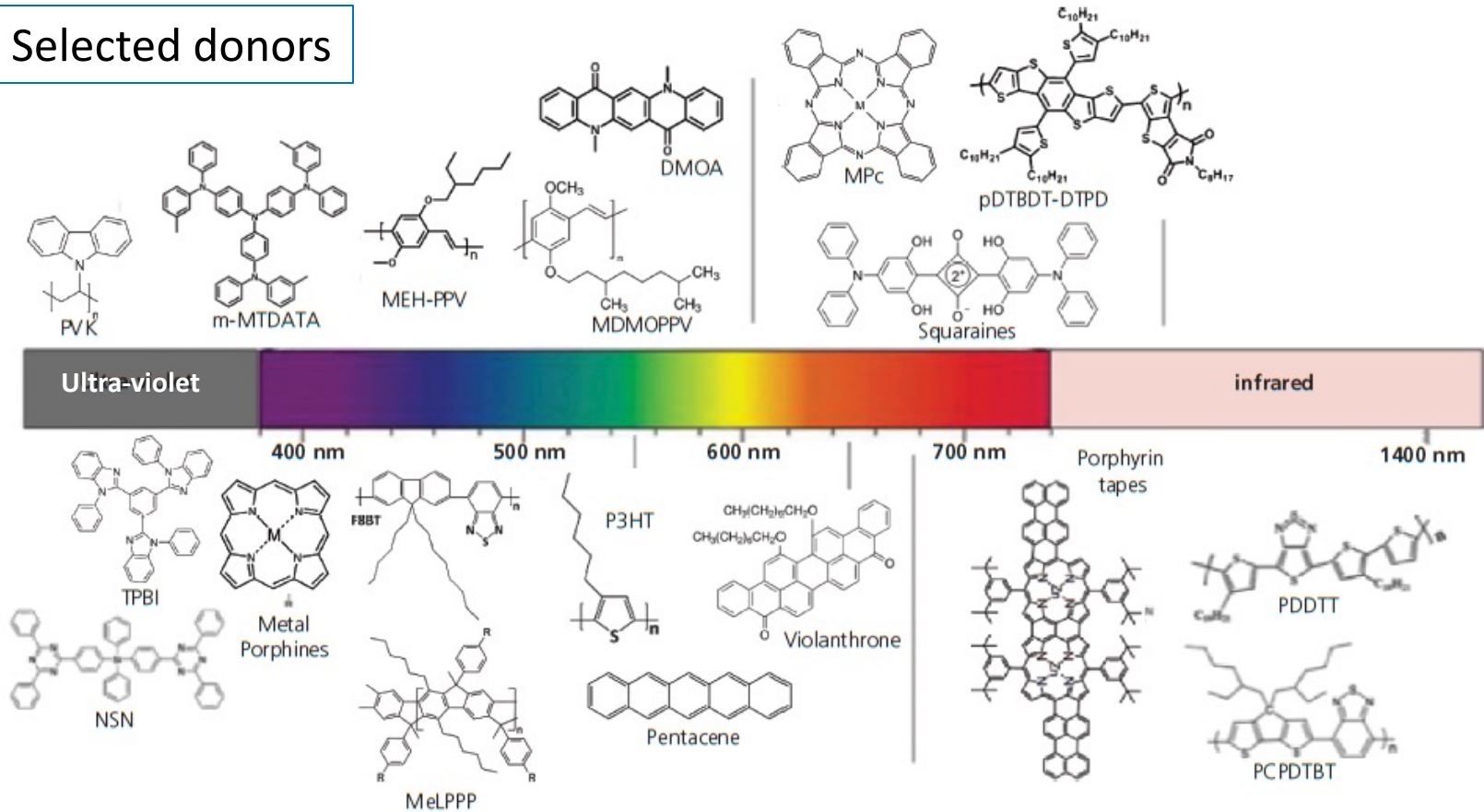
η_p = power conversion efficiency
FF = fill factor

Tang, Applied Phys. Lett., (1986) **48**, 183.

Photodetector Materials

- Good materials absorb in the region of interest
- Morphology promotes exciton diffusion and charge conduction (high mobility)

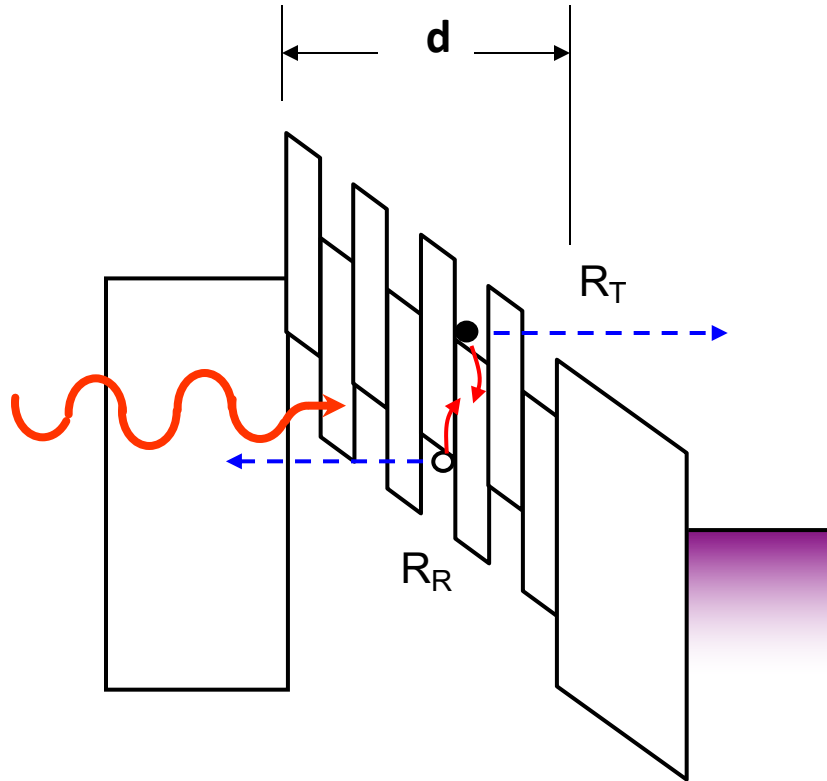
Selected donors



Generally, donors employ fullerene acceptors in OPDs



High Bandwidth Multilayer Photodetectors



Place all D/A junctions within L_D of absorption site

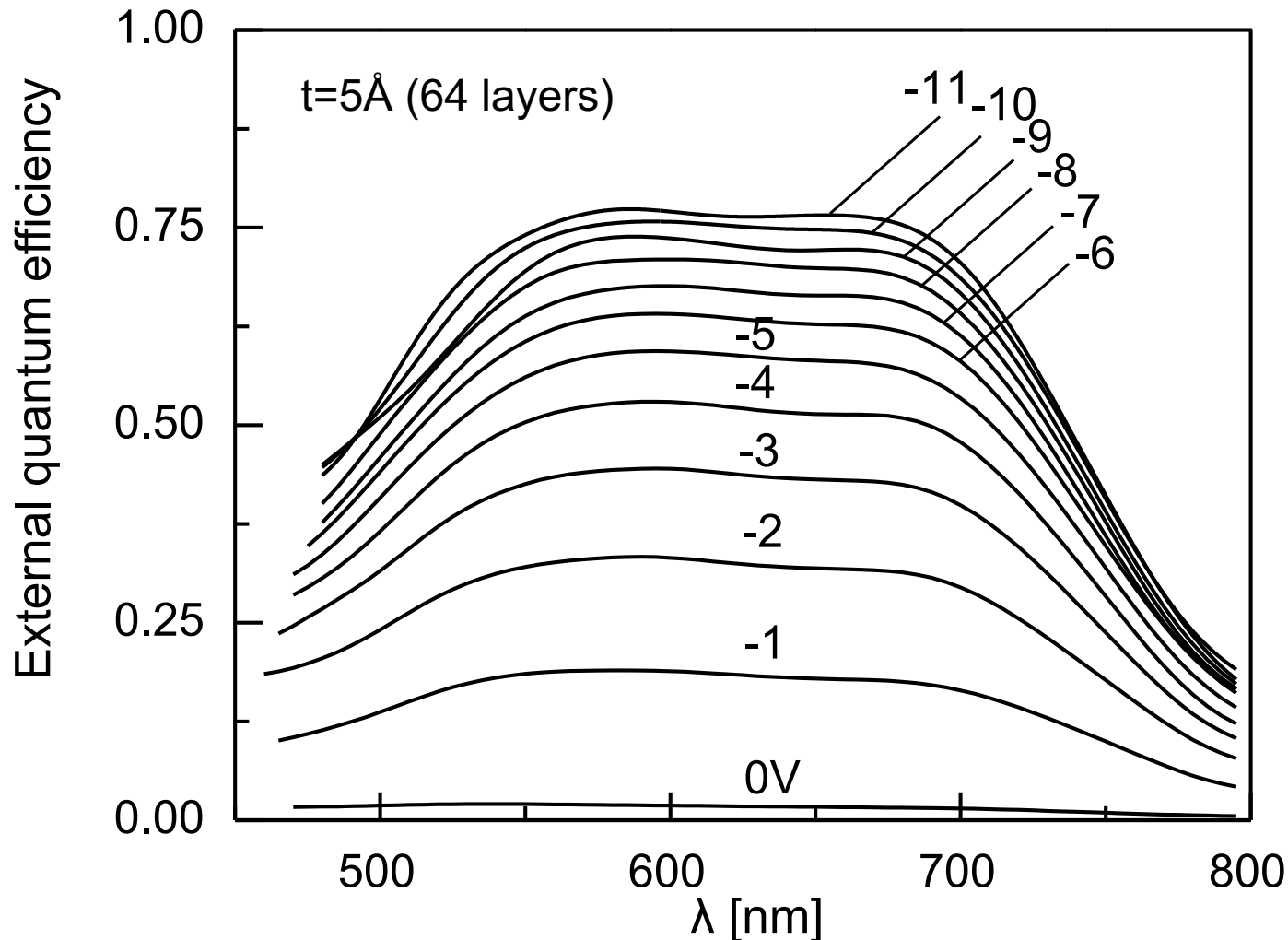
Stack layers until total thickness $d \sim 1/\alpha$

Apply voltage to sweep charge out of potential wells

Bandwidth due to transit time across d .

Spectral + Voltage Dependence of the EQE

- Sensitive to visible + NIR wavelengths
- Strong dependence on bias: EQE ~ 75% @ -10V



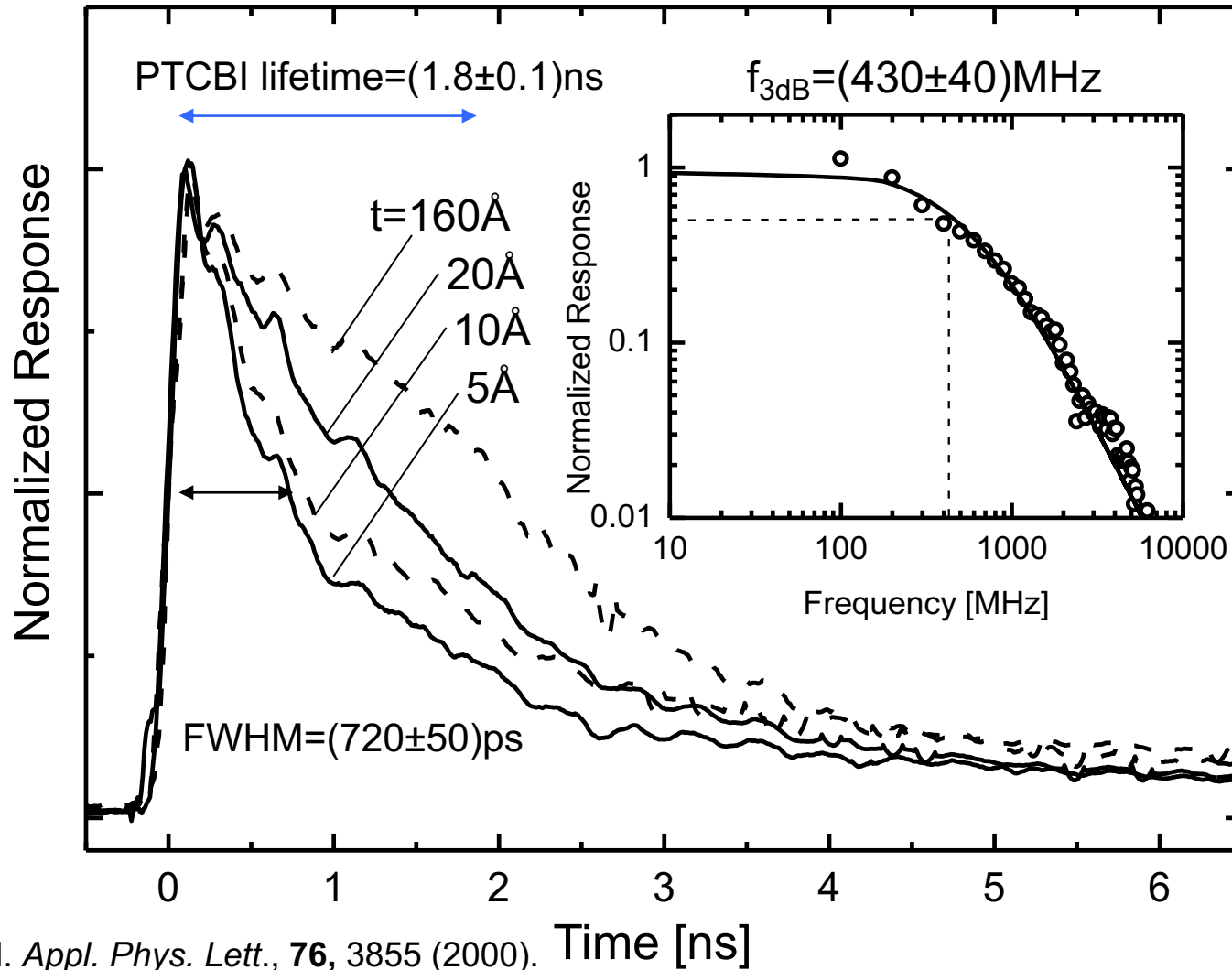
Peumans, et al. *Appl. Phys. Lett.*, **76**, 3855 (2000).

Response Time

Thinner individual layers makes faster devices due to a reduced exciton lifetime

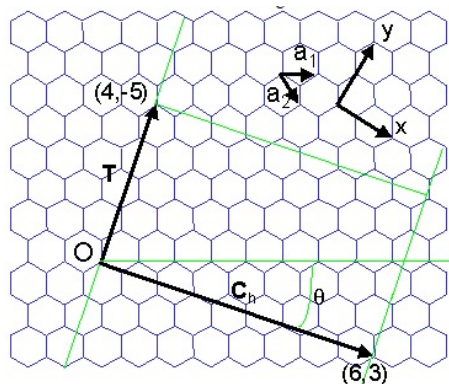
100 μm diameter, -9V, 1.4ps excitation @ 670nm at $(1.0 \pm 0.3) \text{W}/\text{cm}^2$.

Estimated carrier velocities: $v = d/\tau = (1.1 \pm 0.1) \times 10^4 \text{ cm}/\text{s}$

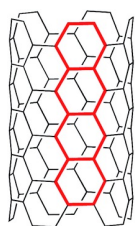


Long wavelength Detectors

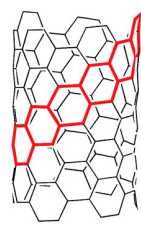
Carbon Nanotubes Can Stretch Detection to NIR



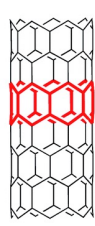
$$C_h = na_1 + ma_2$$



[5,5] CNT
Armchair
Metallic



[7,5] CNT
Chiral
Semiconducting



[7,0] CNT
Zigzag
Semiconducting

Chirality determines if CNT is metallic, semiconducting or insulating

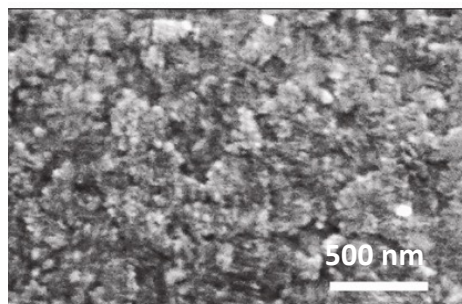
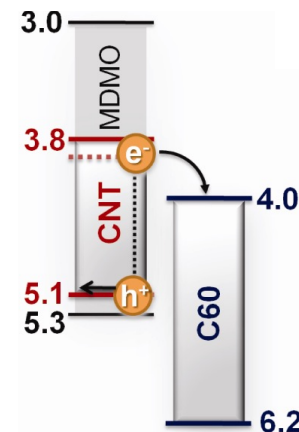
$n = m$: Metallic

$n - m = 3i$ (i integer), $n \neq m$, $nm \neq 0$: semimetal

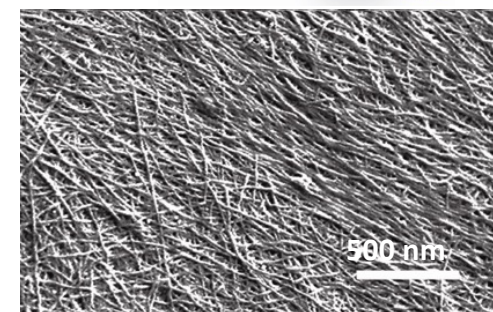
otherwise: semiconductor

Organic/CNT Detector

| | |
|---------------------------|------------|
| Ag | |
| BCP | 10 nm |
| C60 | 100 nm |
| C60 SnPc:C60 | 0 10 nm |
| MDMO-PPV:CNTs P3HT:CNTs | 14 45 nm |
| ITO/glass | |



CNT:MDMO-PPV composite

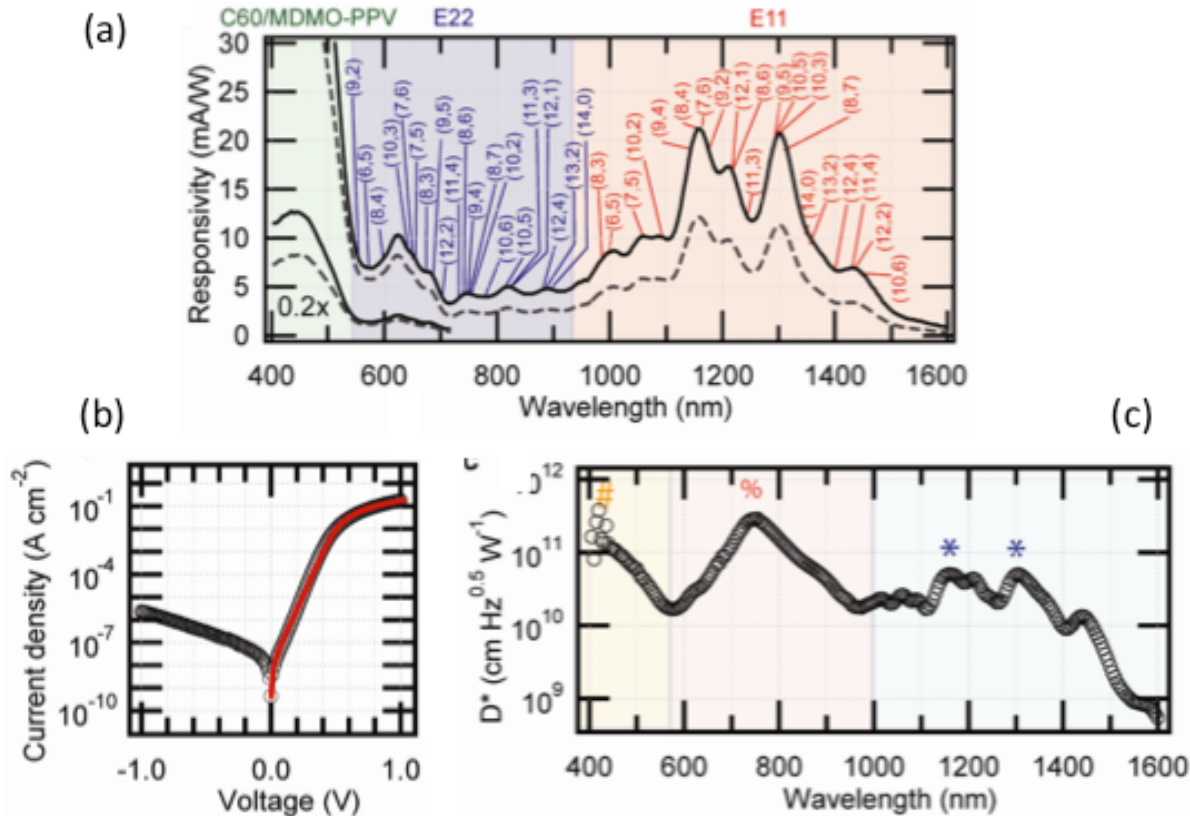


Mat of bare CNT

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Long wavelength Detectors

Single Walled Nanotubes Wrapped in Polymer



Responsivity and Specific Detectivity:

$$\mathcal{R} = \frac{j_{ph} A}{P_{inc}} = qgn_{ext} \left(\frac{\lambda}{hc} \right) \quad [A/W] \quad D^* = \frac{\sqrt{A\Delta f}}{NEP} = \mathcal{R} \sqrt{\frac{A\Delta f}{i_n^2}}$$

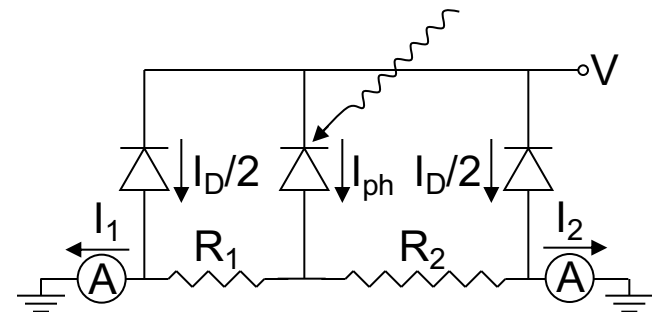
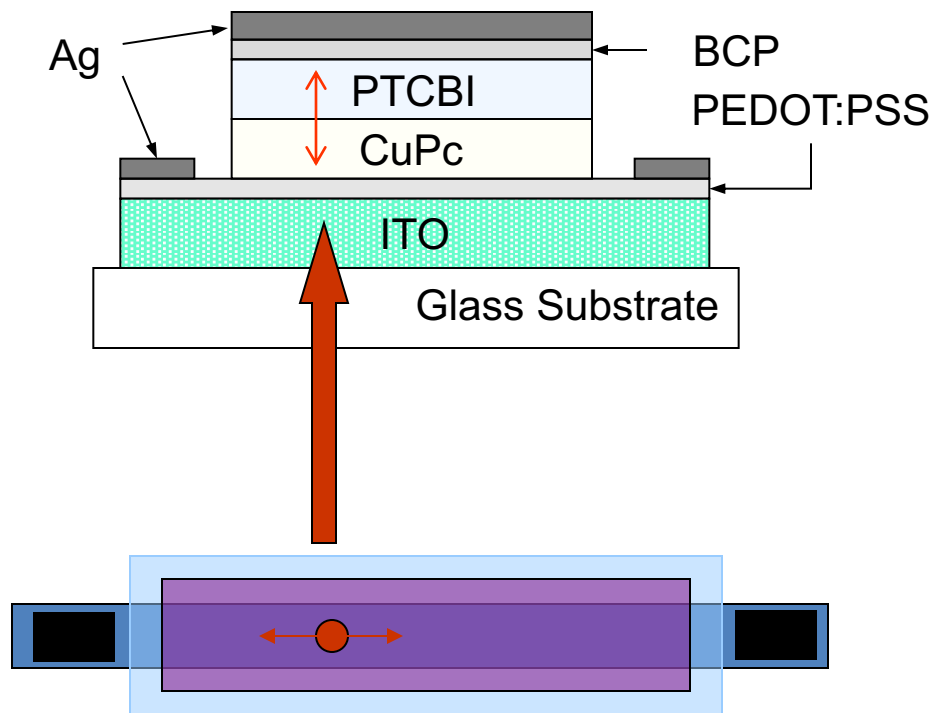
[cm-Hz^{1/2}/W]

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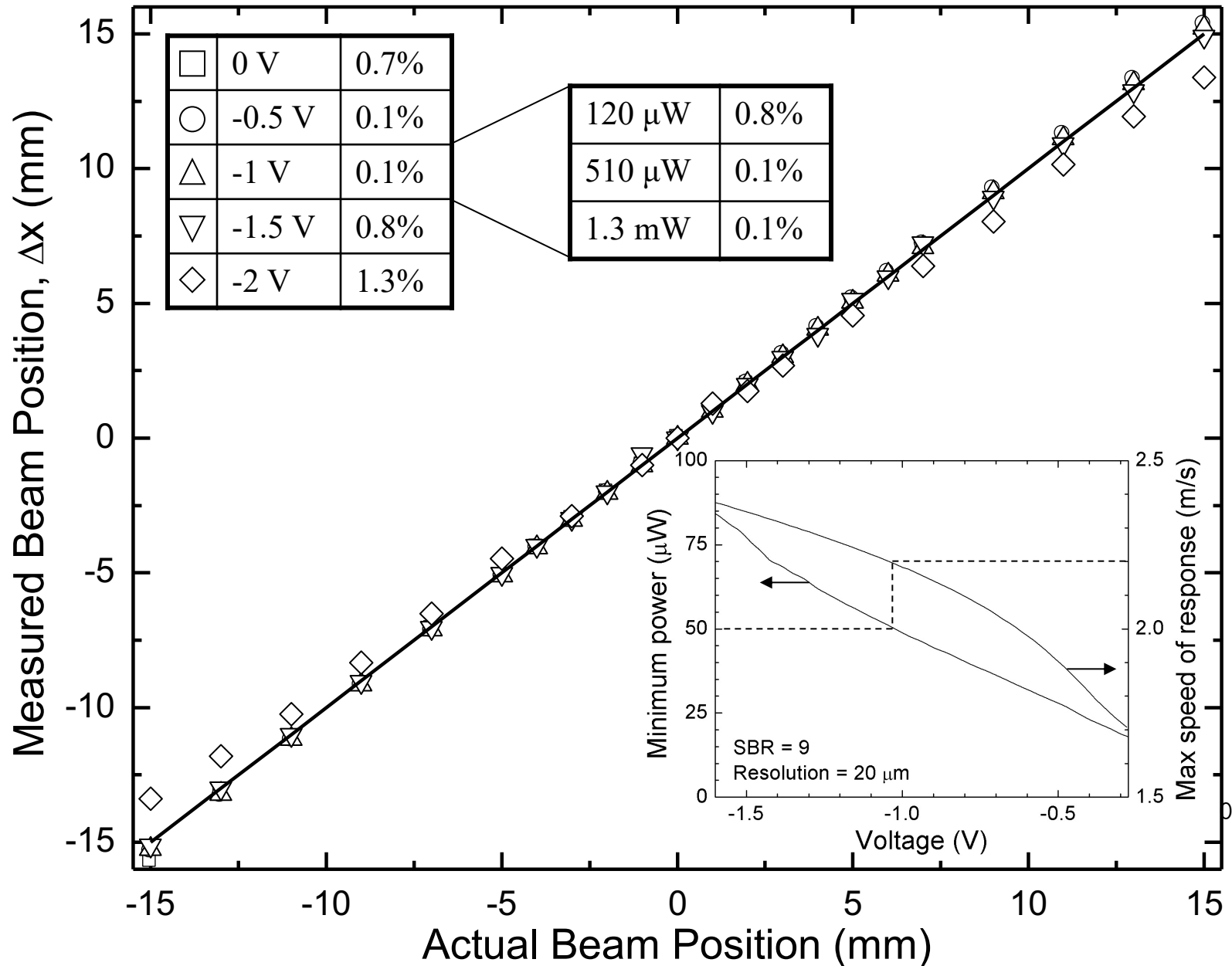
M. S. Arnold, et al., *Nano Letters*, **9**, 3354, 2009.

Position Sensitive Detectors

- Mechanism of operation
 - Extended junction transports charge vertically (no current spreading)
 - Current divided by *linear* resistance of ITO strip



Position Detection Characteristics



Applications of PSDs

- Machine vision
 - Part location and positioning
 - Robot servo feedback
 - 2D possible
- Lab bench positioning
- Free space communication

